

Announcement:

- A few extra positions for enrollment.
- send your (name, ID) to me.
- has to be through my special permission.

$$M \cdot M^T = I$$

Examples of Orthogonal Matrices

① $I = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$ (trivial example)

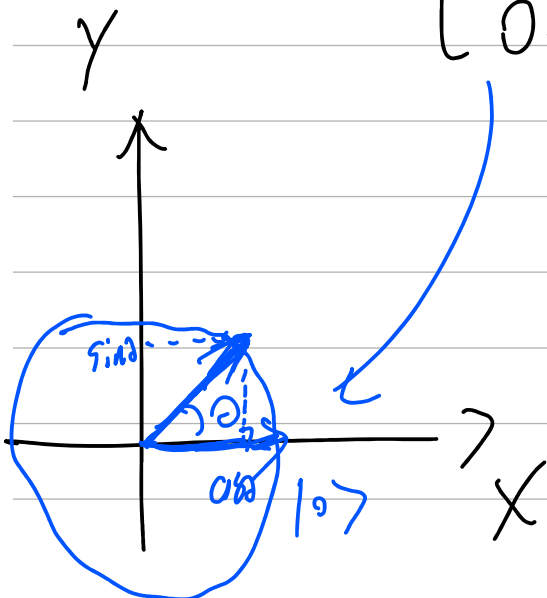
② $R_\theta = \begin{bmatrix} \cos\theta & -\sin\theta \\ \sin\theta & \cos\theta \end{bmatrix} \quad \theta \in [0, 2\pi]$

(counter clock-wise rotation)

$R(-\theta) = ?$

$$R_\theta \cdot \begin{bmatrix} 1 \\ 0 \end{bmatrix} = \begin{bmatrix} \cos\theta & -\sin\theta \\ \sin\theta & \cos\theta \end{bmatrix} \begin{bmatrix} 1 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 1 \\ 0 \end{bmatrix} = \begin{bmatrix} \cos\theta \\ \sin\theta \end{bmatrix}$$



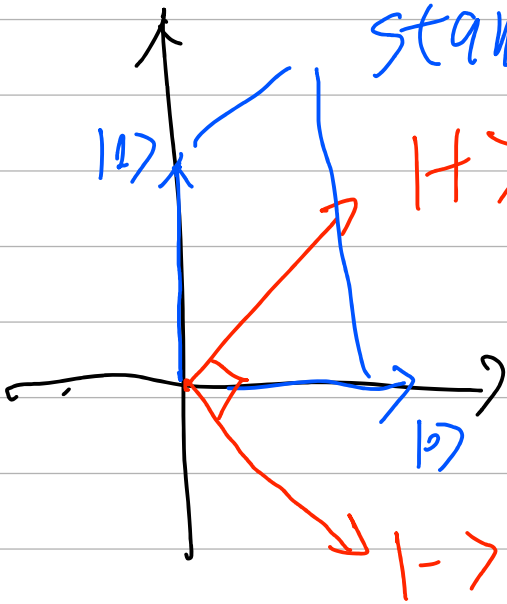
③

$$H = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} \quad (\text{Hadamard gate})$$

$$H \cdot |0\rangle = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} \begin{bmatrix} 1 \\ 0 \end{bmatrix} = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 \\ 1 \end{bmatrix} = |+\rangle$$

$$H \cdot |1\rangle = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} \begin{bmatrix} 0 \\ 1 \end{bmatrix} = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 \\ -1 \end{bmatrix} = |-\rangle$$

- Hadamard basis, computational basis
standard / orthogonal



Hadamard basis

$$\sigma_z |+\rangle = |-\rangle$$

$$\sigma_z |-\rangle = |+\rangle$$

$$\textcircled{3} \quad \sigma_x = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \quad (\text{Pauli } X)$$

$$\sigma_x |0\rangle = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \begin{bmatrix} 1 \\ 0 \end{bmatrix} = \begin{bmatrix} 0 \\ 1 \end{bmatrix} = |1\rangle$$

$$\sigma_x |1\rangle = \dots = |0\rangle$$

superposition

$$\sigma_x (\alpha|0\rangle + \beta|1\rangle) = ? \quad \alpha|1\rangle + \beta|0\rangle$$

NOT-gate in classical computing.

$$\textcircled{4} \quad \sigma_z = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \quad (\text{Pauli } Z) \quad \text{phase}$$

$$\sigma_z \cdot |0\rangle = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \begin{bmatrix} 1 \\ 0 \end{bmatrix} = \begin{bmatrix} 1 \\ 0 \end{bmatrix} = |0\rangle$$

$$\sigma_z \cdot |1\rangle = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \begin{bmatrix} 0 \\ 1 \end{bmatrix} = \begin{bmatrix} 0 \\ -1 \end{bmatrix} = -|1\rangle$$

$$\sigma_z \cdot (\alpha|0\rangle + \beta|1\rangle) = ?$$

$$\sigma_z \cdot |+\rangle = ?$$

(Pauli Y?)

Postulate 4: Null.

It's about how multiple qubits compose into a whole system

Summary:

Postulate 1: A quantum register encode a "unit circle" $| \psi \rangle = \alpha | 0 \rangle + \beta | 1 \rangle$

Postulate 2: When we "measure" a qubit $| \psi \rangle = \alpha | 0 \rangle + \beta | 1 \rangle$, it "collapses" to $| 0 \rangle$ with probability α^2 , and "collapses" to $| 1 \rangle$ with probability β^2 .

Postulate 3: Evolution of a single qubit must be a 2×2 orthogonal matrix multiplied on the left side of the qubit
I.e. $| \phi \rangle = M \cdot | \psi \rangle$, where M is a 2×2 orthogonal matrix

Applications of our over-simplified one-qubit quantum system:

- Truly random number generator
- Elitzur-Vaidman Bomb
- BB84 Quantum Key Exchange / Distribution (QKD)

Truly random number generator

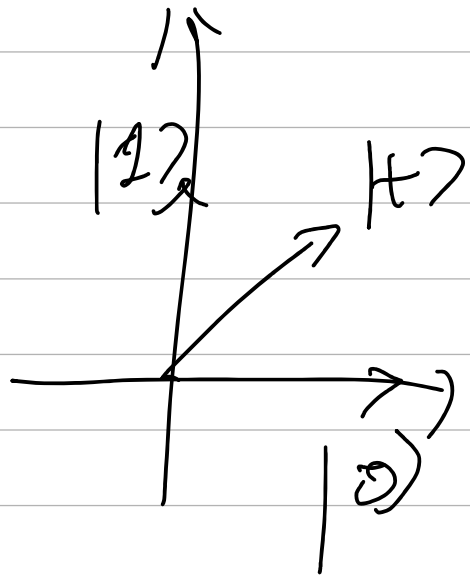
- "True" randomness is expensive,
- You can never be sure about if a source of randomness is truly random in classical physics. (Past Choices.)
 - Sunspot activities / solar phenomenon (solar flares, blackspot bursts)

(highly complex and chaotic processes)

Problematic!

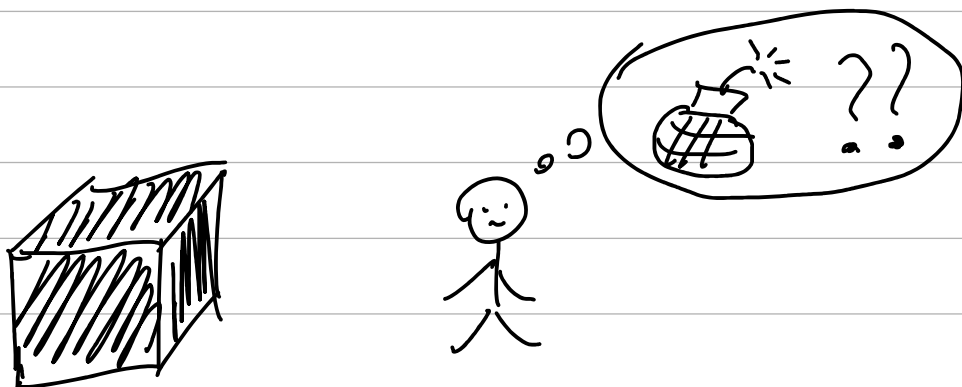
$$H \cdot |0\rangle = |+\rangle = \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle)$$

$$\left[\begin{array}{l} \alpha |0\rangle + \beta |1\rangle \\ \alpha^2 \downarrow \quad \beta^2 \downarrow \end{array} \right]$$



Elitzur-Vaidman Bomb: ^{Shor} 1994

(First proposed in the early 1990's)



If there is a bomb inside, it will immediately explode once the box is opened.

We can solve the problem safely by using a "quantum interrogation"

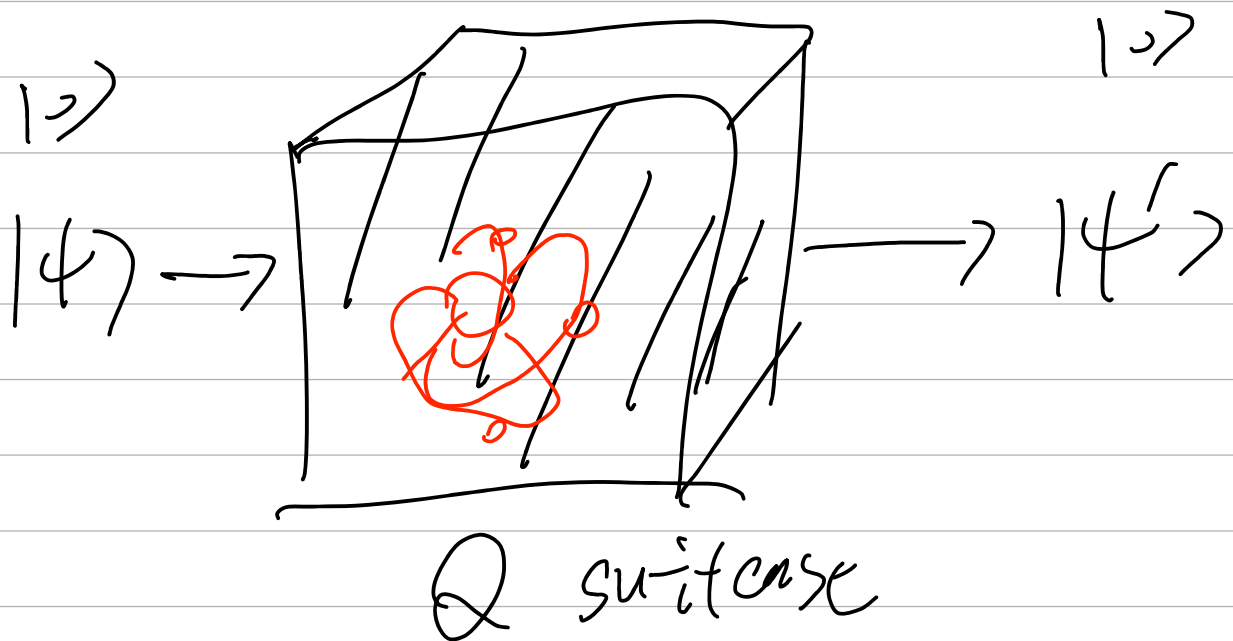
0 — not open } decision.
1 — open

Derivation:

Q. Version

Q. decision qubit

$$|4\rangle = \alpha|0\rangle + \beta|1\rangle$$



If bomb: measure $|4\rangle$ in $\{|0\rangle, |1\rangle\}$

- outcome: $|0\rangle \rightarrow$ nothing. ✓

- outcome: $|1\rangle \rightarrow$ explode X

No bomb:

return $|4\rangle$ back to you, untouched.

R_ϵ . $\epsilon \in (0, 1)$

Pick a number $T = \frac{\pi}{2\epsilon}$

Init $|\psi_0\rangle = |0\rangle$

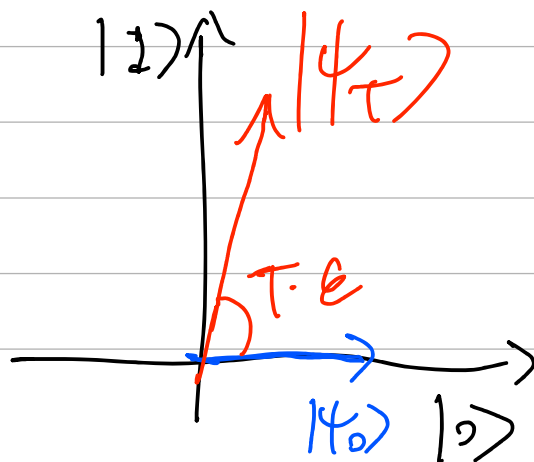
For i in $\{1, 2, \dots, T\}$:

[- Update $|\psi_i\rangle = R_\epsilon |\psi_{i-1}\rangle$
- Query using $|\psi_i\rangle$

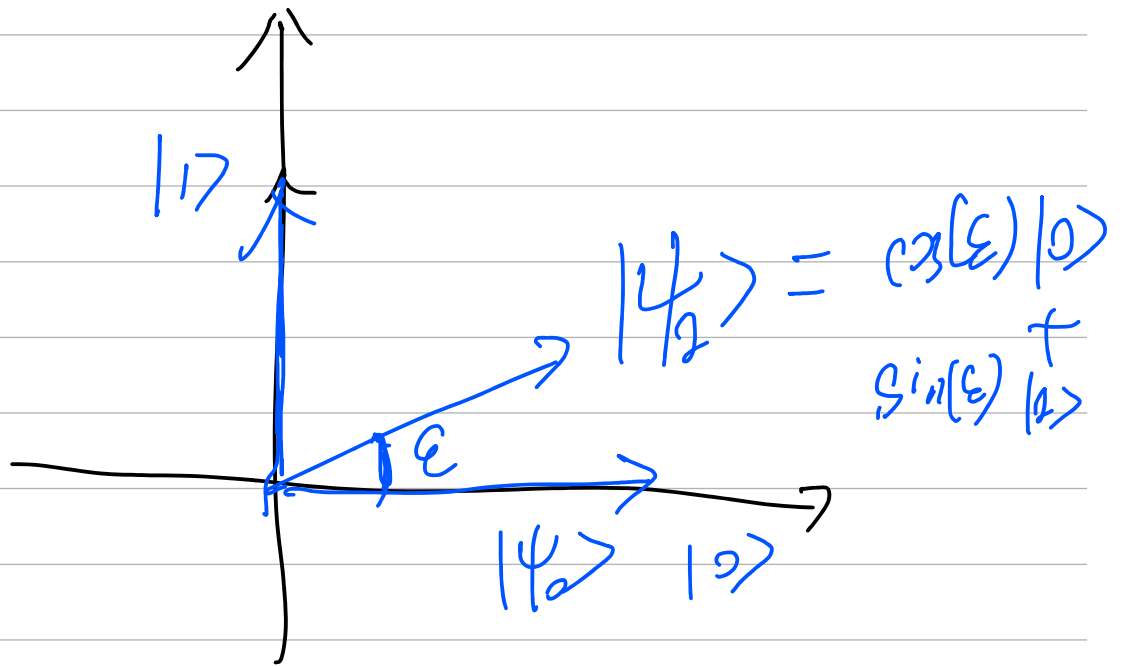
At the end:
measure $|\psi_T\rangle$
- $\{|0\rangle, |1\rangle\}$
comp / stand basis

$|1\rangle$ — no bomb
 $|0\rangle$ — bomb

No bomb.



bomb



{ Safe: w.p. $\cos^2(\epsilon)$

{ Explode: w.p. $\sin^2(\epsilon) \approx \epsilon^2$

$\epsilon = 0.00000001$ $\left[\begin{array}{l} \sin(x) \sim x \\ \text{if } x \text{ small} \end{array} \right.$

Assume safe
for T times, $|\psi_T\rangle = |0\rangle$

Safe_T

$\Pr[\text{Safe}_T] \geq 1 - T \cdot \sin^2 \epsilon$
 $= 1 - T \cdot \epsilon^2$

Union bound

$$T = \frac{\pi}{2\varepsilon}$$

$$1 - T \cdot \varepsilon^2 = 1 - \underbrace{\frac{\pi}{2} \varepsilon}$$

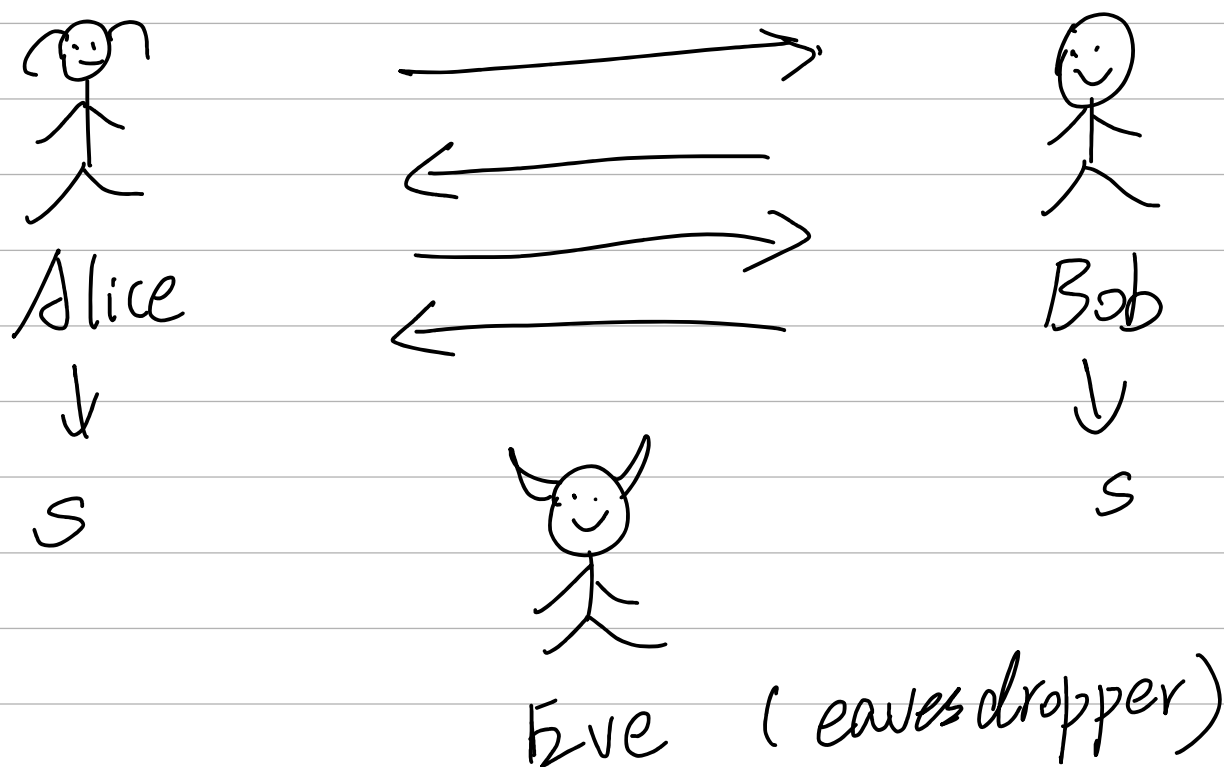
Risk tolerance bound

$$\varepsilon = \frac{0.0000001}{\pi/2}$$

$$\Pr(\text{safe}) \geq 1 - 0.0000001$$

Year 1984

BB84 Quantum Key Distribution (QKD)



Classical Solution:

Public-key Encryption (PKE)

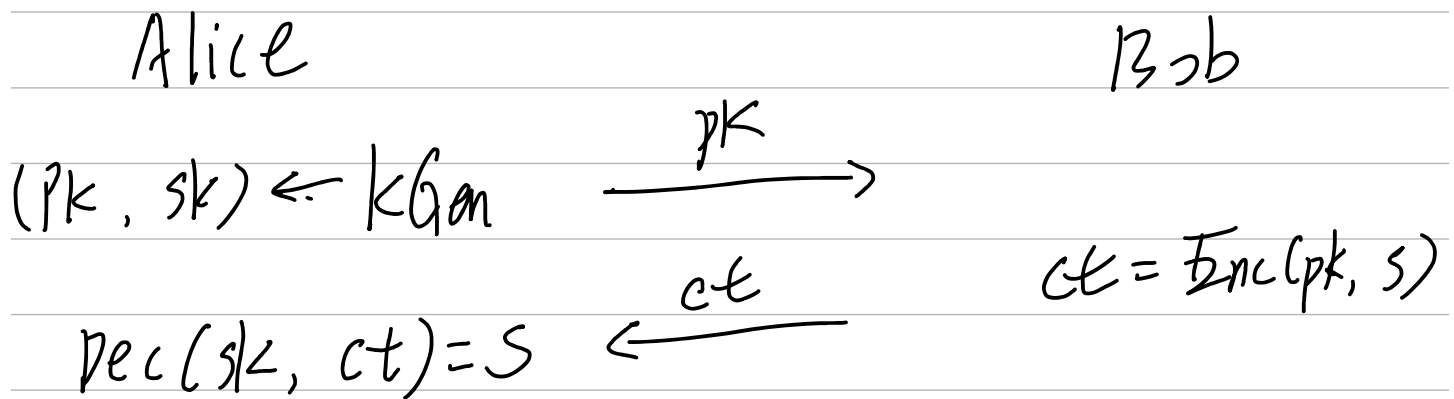
$$\left\{ \begin{array}{l} \text{KGen} \rightarrow (pk, sk) \\ \text{Enc}(pk, m) \rightarrow ct \\ \text{Dec}(sk, ct) \rightarrow m \end{array} \right.$$

SRSA
DH
Factoring ↙
Dislog

Two requirements for PKE:

- Correctness,

- Security;



"Problems:"

- PKE requires "hardness assumptions"

RSA, Factoring.

- Classically, PKE is necessary for this task,

\Rightarrow "assumptions" are necessary for classical key distribution,

The BB84 Solution:

Alice

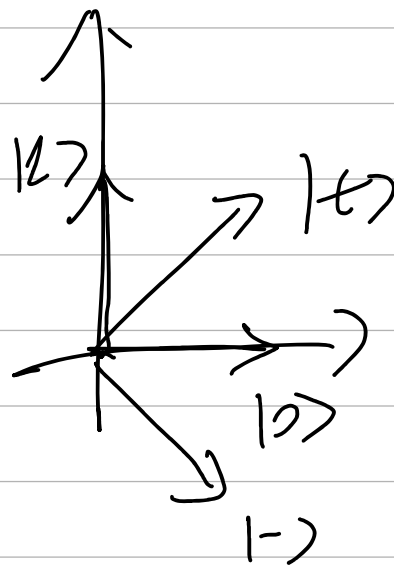
Bob

For $i = 1$ to n .

$$x_i \leftarrow \{0, 1\}$$

$$y_i \leftarrow \{0, 1\}$$

$$|\psi_i\rangle = \begin{cases} |0\rangle & x_i = 0, y_i = 0 \\ |1\rangle & x_i = 0, y_i = 1 \\ |+\rangle & x_i = 1, y_i = 0 \\ |-\rangle & x_i = 1, y_i = 1 \end{cases}$$



$$\underline{|\psi_1\rangle \dots |\psi_n\rangle}$$

For $i = 1$ to n

$$x'_i \leftarrow \{0, 1\}$$

measure $|\psi_i\rangle$

$$\begin{aligned} & \{x_i\}_{i=2}^n \\ & \leftarrow \{x'_i\} \end{aligned}$$

if $x'_i = 0$,

in $\{|0\rangle, |1\rangle\}$ basis

in $\{|+\rangle, |-\rangle\}$ basis

if $x'_i = 1$.

Outcome: y'_i

set s be
concatenation of

x_j 's where

$$j \in \{j : x_j = x_j'\}$$

set s be
concatenation of

x_j 's where

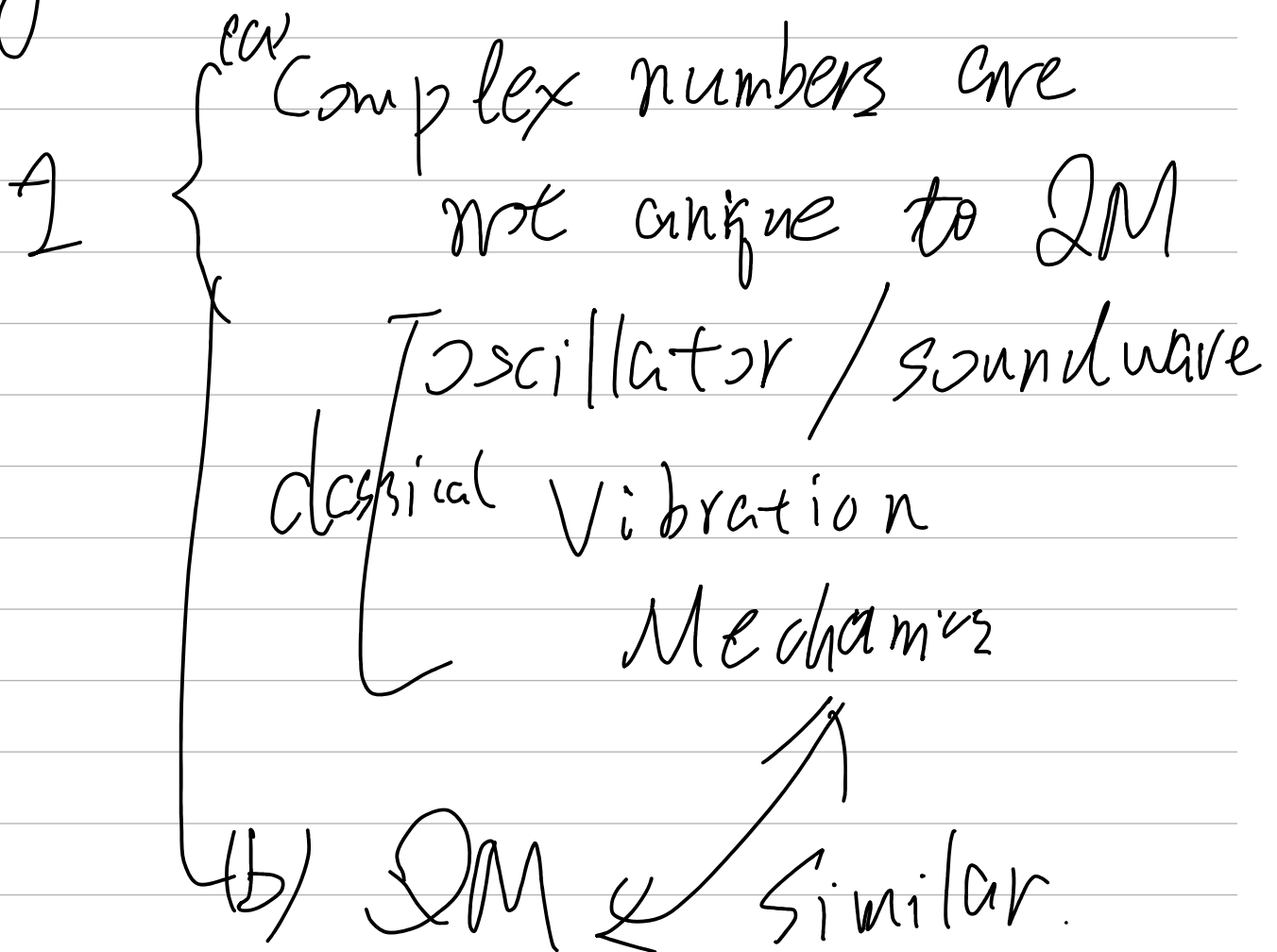
$$j \in \{j : x_j = x_j'\}$$

Why complex numbers in QM?

$$\alpha |2\rangle + \beta |0\rangle$$

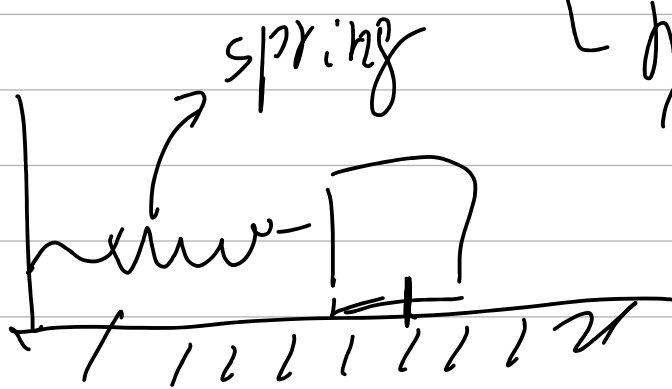
short - answer: Law of QM

long - answer:



2. No. QM needs complex numbers in an essential way.

{ hyp 1: Real-QM
hyp 2: Comp-QM



Simple harmonic Oscillator.

(^力 _い) ^い _い 振動)

$$\left\{ \begin{array}{l} F = -kX \rightarrow \text{position} \\ \downarrow \\ \text{hook parameter} \\ F = m\vec{a} \end{array} \right.$$

$$\left(\frac{d^2}{dt^2} x \right) + \left[\frac{k}{m} \right] x = 0$$

$$\left[i \hbar \frac{d}{dt} |\psi\rangle = \hat{H} |\psi\rangle \right]$$

$$x(t) = C \cdot \cos(\omega t + \varphi)$$

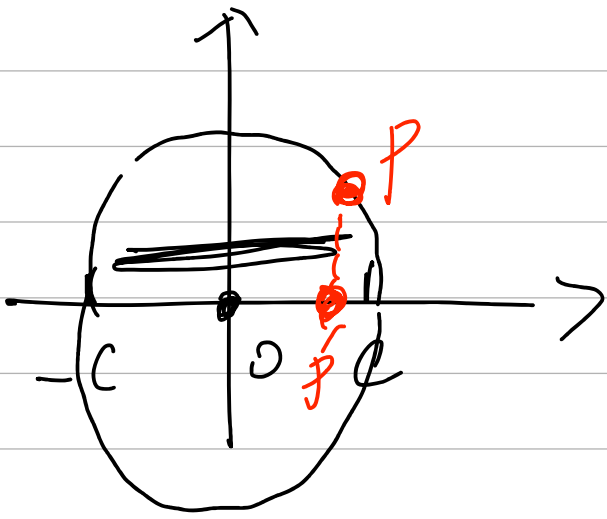
full solution

angular
frequency.

$$x(t) = C \cdot e^{i(\omega t + \varphi)}$$

$$= C \cdot \left[\cos(\omega t + \varphi) + \right.$$

$$\left. i \sin(\omega t + \varphi) \right]$$



$$\int_0^2 \cos^2(x) dx$$

Full-fledged Postulates of (complex-number) QM

Postulate 1: An isolated quantum system is completely described by its vector of state, which is a unit vector in a Hilbert space.

vector space (Real or complex)
inner-product space
complete.